

## ABSTRACTS OF CONTRIBUTED PAPERS

TEMPORAL EVOLUTION OF THE  $H_2$   
ABUNDANCE IN THE INTERSTELLAR MEDIUM

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We calculate the abundance of molecular hydrogen ( $H_2$ ) for a wide range of interstellar conditions, including the following cases: *i*) constant density and temperature, *ii*) constant density and variable temperature (isochoric case), and *iii*) constant pressure (isobaric case). Starting with a diffuse atomic medium, the time evolution of the  $H_2$  abundance can be followed with  $dn_{H_2}/dt = R_f - R_{ph,d} - R_d$ , where  $R_f$  is the formation rate of  $H_2$  on interstellar grains,  $R_{ph,d}$  is the photodissociation rate, and  $R_d$  is the collisional dissociation rate. The formation and collisional dissociation rates depend on the gas temperature and density, and a self-consistent treatment of the problem should include the effects of  $H_2$  cooling. Photodissociation of interstellar  $H_2$  is initiated by absorption of discrete lines of the Lyman and Werner series, and within the clouds the general interstellar radiation field is quickly attenuated at those wavelengths. As the optical depth increases, the photodissociation rate decreases and becomes negligible at large optical depths ( $\tau_v > 1$ ). Also, in the absence of local energy sources, at large optical depths the collisional dissociation rate becomes negligible (i.e.,  $R_d$  tends toward zero at temperatures below  $T \sim 2000^\circ$  K). For all three cases considered, we find that the time required to convert atomic to molecular hydrogen is longer than the timescale for a number of cloud formation mechanisms. For the isochoric case we find that the molecular formation timescale can be written as a power-law of the total density,  $t_{for} = A_0(n/1 \text{ cm}^{-3})^\alpha$  yr, with  $A_0 = 1.06 \times 10^{10}$  and  $\alpha = -1.02$ . Similarly,

for the isobaric case, the formation timescale can be written as  $t_{for} = B_0(P/10^{-12} \text{ dyn cm}^{-2})^\beta$  yr, with  $B_0 = 1.7 \times 10^8$  and  $\beta = -0.88$ .

IMAGERY AND SPECTROSCOPY OF THE  
REGION AROUND HH80–81

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Preliminary results of  $H\alpha$  imaging and low dispersion spectrophotometry of the HH80–81 region are presented. Data were obtained at the Observatorio Astronómico Nacional, San Pedro Mártir B.C., México. The spectral range from 4200 to 7200 Å was covered with a resolution of 8 Å. Spectrophotometric results for HH80 show that it is a high excitation object; our data are similar, within the errors, to data reported by Reipurth & Graham (1988, A&A, 202, 219). We derived an electron density  $N_e([S \text{ II}]) = 3000 \text{ cm}^{-3}$ , so HH80 has an intermediate density compared to other HH objects.

Continuum-subtracted  $H\alpha$  images clearly show the HH80 and 81 objects. A stellar object with strong  $H\alpha$  emission, located to the southwest of HH80, was also detected. This stellar source has not been reported previously. The spectrum of this object shows a stellar continuum with wide Balmer absorption lines;  $H\beta$  and  $H\gamma$  have  $EW \sim 15$  Å. The  $H\alpha$  profile has two components: the photospheric absorption line and an intense, narrow (unresolved) emission line. These spectral characteristics are similar to those of some Ae/Be Herbig stars (Cohen & Kuhl 1979, ApJS, 41, 743). From the calibrated flux at  $\lambda 5500$  and  $\lambda 4400$  we derived apparent magnitudes  $m_V = 12.3$  and  $m_B = 12.5$ , respectively. By assuming a distance of 1.7 kpc to the HH80–81 complex (Martí, Rodríguez, & Reipurth 1993, ApJ, 416, 208), we found an absolute visual magnitude  $M_V = -1.0$  for the stellar object, which is consistent with a spectral classification of B9e.